Structural controls on the location of Longqi hydrothermal activity at inside corner revealed by the stress state

Ming Chen
Chunhui Tao (taochunhuimail@163.com)
Key Laboratory of Submarine Geosciences, SOA, Second Institute of Oceanography

Yunlong Liu
Hanchuang Wang
Weifang Yang
Kaiying Chen
Mingxu Wang
Dong Liu
Qianyu Li
Dong Chen

Article

Keywords:

Posted Date: May 26th, 2022

DOI: https://doi.org/10.21203/rs.3.rs-1642852/v1

License: This work is licensed under a Creative Commons Attribution 4.0 International License.
Read Full License
Abstract

The increased hydrothermal activity at inside corners is interpreted to be controlled by the high-permeability detachment faults at mid-ocean ridges. Paradoxically, almost none of hydrothermal activity is confirmed to be located at the detachment terminations, where the permeability is theoretically maximal. Here, we use a numerical model to investigate the stress state of the Longqi hydrothermal field at inside corner, and then explore the structural controls on the hydrothermal activity. The results show that inside corner is undergoing regional stress rotation, deviating from the spreading direction by more than 45°. Shear strain resulting from stress rotation enhances the regional permeability of the shallow crust and therefore might control the location of the hydrothermal site. Our model provides a potential mechanism for hydrothermal activity in similar geological settings where the shear strain enhances regional permeability and facilitates hydrothermal circulation in the hanging wall, and thus reconciles the disparity between prevailing models and field observations.

Introduction

It has long been recognized that on the slow and ultraslow spreading Mid-Ocean Ridges (MORs), the axes are typically divided into distinct accretionary segments by discontinuities, creating fundamental segmentations of the lithosphere plates at a range of scales. Inside corners formed either at non-transform discontinuities (NTDs) or at Ridge-Transform Intersections (RTIs) are often associated with increased heat flow and hydrothermal activity. A possible explanation for hydrothermal activity in such settings is that hydrothermal circulation and the generation of the large mineral deposits are dominated by the long-lived, high-permeability detachment fault, which penetrates into Brittle-Ductile Transition (BDT, ~650°C) and allows the hydrothermal fluids to tap heat from a deep reservoir. Up to now, all the Oceanic Core Complexes (OCCs) that have been explored on the Mid-Atlantic Ridge (MAR) are known to either host active hydrothermal vents or show extinct hydrothermal deposits. Paradoxically, none of the hydrothermal sites is located at the termination of the detachment, where the permeability is reasonably maximal. In addition, hydrothermal activity is not restricted to the relatively warm segment centers, where the thermal flux is reasoned to be at a maximum. On the contrary, it can equally be located at relatively cold segment ends and discontinuities, which could not be fully explained by previous models.

Constructions of hydrothermal vents require regional permeable pathways and the upwelling of hydrothermal fluids focused at this site. Observations of the Trans-Atlantic Geotraverse (TAG) hydrothermal field, which is located at the hanging wall of a detachment fault, showed that multiple phases of extensional faulting and oblique fractures occurred within the hydrothermal field and the localization of hydrothermal activity seems to be strongly controlled by the permeable pathways. The cause of these oblique faults is most likely not associated with ridge-parallel faults as they deviate from the spreading direction by more than 40°. Up to now, what still remains problematic is the origin of these oblique faults that sustain the upwelling of hydrothermal fluids.
In this paper, we use a numerical model to investigate the present-day stress state of the Longqi hydrothermal field on the ultraslow spreading Southwest Indian Ridge (SWIR). The active Longqi-1 hydrothermal vent (Fig. 1b) is located at the hanging wall of a detachment fault at inside corner and therefore provides an opportunity to explore the possible mechanism for hydrothermal activity in such settings. For comparisons, we estimate stress state of TAG hydrothermal field on the MAR, which is located at a similar geological setting with Longqi field. In contrast to previous models\textsuperscript{19–21}, bathymetry and hydrostatic pressure are coupled in our numerical model in order to investigate the structural controls on the hydrothermal activity, and then understand how a favorable tectonic and volcanic setting may develop for hydrothermal circulation. We will show that shear strain resulting from local stress rotation enhances the regional permeability of the shallow crust and controls the location of the hydrothermal activity.

**Geological Setting**

The Longqi segment lies at the center of the SWIR (49.3–49.9°E) between Segment 28 and 29, where a ~19.5-km-long NTD is located in between (Fig. 1b). The half-spreading rate is 5.2 mm/y in the north flank and 8.7 mm/y in the south flank, respectively\textsuperscript{22}. The average water depth is shallower in the south flank (< 2000 m) than that in the north flank (> 2000 m). Using the 7.0 km/s isovelocity contour as a reference for the base of the crust, seismic experiments showed that the crustal thickness at inside corner (Fig. 1b) was ~ 2.0–3.0 km while that at outside corner (Fig. 1b) is 4–5 km\textsuperscript{23,24}. The P-wave velocity at insider corner is ~ 0.5-2.0 km/s higher than average P-wave velocity in the shallow oceanic crust\textsuperscript{25}. The fraction of the plate separation rate accommodated by magma emplacement, $M$, is ~ 0.48\textsuperscript{26}. A possible explanation for the observed structural asymmetry therefore is that nearly all the magmatic accretion occurs in the north flank with small-offset faults while seafloor spreading in the south flank is mainly accommodated by the large-offset detachment fault\textsuperscript{23,24}, as suggested by numerical models\textsuperscript{27–29} and field observations\textsuperscript{10,30}.

The first active, high-temperature Longqi-1 hydrothermal vent field on the SWIR was confirmed during Chinese cruise DY115-19 in 2007\textsuperscript{31}. The highest measured temperature of the venting fluid is as high as 350°C\textsuperscript{12}. Near-bottom magnetic data shows a well-defined low crustal magnetization zone\textsuperscript{32}, which is interpreted be a consequence of hydrothermal alteration. The basalt-hosted Longqi-1 hydrothermal vent is located on the hanging wall of the detachment fault at inside corner. Microearthquake study reveals a relatively uniform BDT base along Longqi segment and adjacent NTD, indicating that no focused melting took place beneath this segment\textsuperscript{25}. Likewise, there is no relatively low-velocity domain beneath the Longqi hydrothermal field\textsuperscript{23}. Heat source thus would most likely be the melt zone below the Axial Volcanic Ridge (AVR)\textsuperscript{12}. The detachment fault provides the high-permeability pathway for hydrothermal circulation from upper mantle to the seafloor\textsuperscript{12,23}.
A key point in geomechanical modeling is the reliability of the numerical results in terms of the predicted stress field and mechanical process. Uncertainties are mainly introduced by the model input (e.g. rock properties, boundary conditions and structures). To test the reliability of our model, we compare the present-day stress field predicted by our numerical model with that estimated by model-independent earthquake focal mechanisms. Earthquakes recorded by a sufficient number of stations with a reasonable azimuthal distribution are usually used to estimate focal mechanisms and infer the orientation of the maximum horizontal stress ($S_{Hmax}$). The focal mechanism database used in this study is obtained from the Global Centroid Moment Tensor Project, where more than 150 stations are available to calculate the stress tensor. The focal mechanisms used in this study are shown in Fig. 2 (beachballs). Table 2 illustrates the comparisons of $S_{Hmax}$ orientation inferred from the focal mechanisms and numerical model, respectively. The maximal difference (C) is 10° and the average difference is ~ 5°, suggesting that the orientation of the stress field predicted by numerical model has a good match with the observed. However, it should be noted that the focal mechanisms indicate the stress field at depth of ~ 10 km below seafloor while the numerical results show the present-day stress state of the seafloor.

The orientation and magnitude of the $S_{Hmax}$ predicted by our model are shown in Fig. 2 (black arrows). The magnitude of $S_{Hmax}$ ranges from 15 to 40 MPa and the orientation shows a reasonable correlation with the observed spreading direction. At the inside corner, the $S_{Hmax}$ is rotated by approximate 45° relative to the ridge trend, indicating a strike-slip regime. The magnitude of the shear strain predicted by our model in the study area are shown in Fig. 2b. The magnitude varies from $-3 \times 10^{-4}$ to $3 \times 10^{-4}$, where positive values represent the right-moving shearing process and the negative values represent the left-moving shearing process. Similar with the stress field, the distribution of shear strain between the opposite flanks is strongly asymmetric. Shear strain at inside corner is obviously higher than that at outside corner. The red, yellow and black stars represent the active Longqi-1 hydrothermal vent (red star), inactive Longqi-3 hydrothermal field and hydrothermal anomaly.

**Discussion**

**Stress rotation and oblique faults at inside corners**

Similar to many other segments on the slow and ultraslow spreading ridges, the Longqi segment are characterized by the strong interplay between tectonic and magmatic processes. Shipboard bathymetry shows a strong asymmetry between the north and south flank (Fig. 1b). The north flank is characterized by narrow, northward-dipping, axis-parallel normal faults that extend up to 40 km. The small-offset faults reflect relative rich magma supply. In contrast, the fault heave and space in the south flank are obviously larger and more spatially complex, reflecting reduced magma supply and stronger interplay between volcanic and tectonic process. In the south flank, seafloor spreading is dominated by the long-lived detachment fault capped by corrugated slip surfaces. Faults at the segment centers strike approximately E-W in respond to extension. At the segment end, however, topographic evidence reveals
the presence of three oblique faults in a NE-SW trend (dashed red lines in Fig. 1b). Applying a selection criterion based on a root mean square fault plane uncertainty of $< 35^\circ$ and a mechanism probability of $> 60\%$, Yu et al. (2018) reported two well-constrained earthquake focal mechanisms located $\sim 3$ km west of the Longqi-1 hydrothermal vent. These events indicate almost N-S-trending fault planes and a strong shear component in the lithosphere. Although the origin of these oblique faults is not clearly understood, we believe the mechanism for the localized strike-slip structures must be likely not as same as that for the ridge-parallel faults as they deviate from the general trend of the spreading axis by more than $40^\circ$. A possible explanation for the oblique structures is that the oblique faults are associated with the NTD in the west. Previous studies on the MAR revealed three significant mechanical characters of the NTDs: 1) NTDs are dominated by oblique fault scarps; 2) sustained strike-slip occurs within the NTDs; 3) the geometry of NTDs appears to be unstable. The non-rigid offset represents a rheologically weak plate boundary along which shear stress is accommodated over a broad zone of oblique, normal and strike-slip faulting. Shear deformation at NTDs therefore appears to be accommodated by wide zones ($> 10$ km) of oblique or ridge-parallel faults instead of the narrow zones ($< 5$ km).

The Longqi-1 hydrothermal vent field, located at the inside corner, is mapped in higher resolution (Fig. 3). The bathymetric data were collected by the autonomous underwater vehicle Autonomous Benthic Explorer (ABE) and operated by the Woods Hole Oceanographic Institution (WHOI) during the second leg of the Chinese cruise DY115-19 on board R/V DaYangYiHao in 2007. The depth is $\sim 2650$ m and its summit is a flat ($\sim 15^\circ$) and corrugated surface, interpreted to be an exhumed detachment surface. More than 20 active and inactive vents were reported based on near-bottom Jiaolong human-occupied vehicle and Qianlong II autonomous underwater vehicle. All of these vents are located along the narrow faults (Fig. 3), which are oblique or perpendicular to the spreading axis (Fig. 3f, g) and are restricted within a 0.5-km-wide band near the active vents. We observed a series of $\sim 20$-m-high eastward-dipping faults, which strike NNW (Fig. 3) and are located $\sim 10$ m west of the active vent. The presence of basalt breccias (Fig. 3d) and the oblique faults (Fig. 3g, h) in NE-SW and W-E trend directly observed by deep-tow camera indicates the rotation of the local stress. The distribution of hydrothermal vent sites suggests that hydrothermal fluid should be focused in the shallow crust beneath oblique faults. The location of hydrothermal vents seems to favor an interpretation where the oblique faults enhance the regional permeability of the hanging wall of the detachment fault and further facilitate hydrothermal circulation. Outflow from a conduit with such a narrow linear geometry would probably generate an axisymmetric sulfide deposit, therefore could explain that the hydrothermal vents are located on the both sides of the faults.

**Mechanical asymmetry between inside and outside corner**

We estimated the present-day stress field of Longqi segment through numerical approaches incorporating bathymetry, material properties and the spreading rate. Our results show a strong asymmetry between inside and outside corners in terms of stress (Fig. 2a) and shear strain (Fig. 2b). The $S_{Hmax}$ is consistent with the observed spreading direction at outside corner. By contrast, at inside corner, the $S_{Hmax}$ is oriented approximately NE-SW, deviating from the ridge trend by more than $45^\circ$ (Fig. 2a).
Meanwhile, the shear strain at inside corner is significantly higher than that at outside corner. With the distance from the NTD, the orientation of the $S_{Hmax}$ is gradually transformed from a predominantly NE-SW direction to a predominantly N–S direction. At the same time, the shear strain decreased from the segment end to segment center. The deviation of stress orientation and related shear strain are most likely localized because the shear stress regime is restricted at inside corner rather than in the entire south flank. The mechanical effects of the NTD are inferred to be the key drivers of the origin of stress rotation and shear strain at inside corner$^{21}$. In the shear stress regime, due to relatively weak and unstable lithospheric edges along the NTD relative to adjacent segments, stress rotation and inelastic deformation most likely initiates at this weaken zone in response to far-filed extension.

Previous field observations and numerical studies have revealed that the asymmetric spreading was controlled by the fraction of magmatic accretion$^{27,29,44}$. In these models where $M$ is $\sim 0.5$, extension at inside corner is mainly accommodated by a long-lived detachment fault. Our results indicate that NTDs play an important role in the formation of the asymmetric volcanic and tectonic fabrics. The non-rigid NTDs show a weaker resistance to the ridge-normal stress, and thus the locations and structural geometry can adjust relatively easily$^1$. The imposition of a shear couple across the NTD results in the rotation of local stress field and the resultant oblique tectonic fabrics. Oblique faults at inside corner have also been recognized from bathymetric data on the MAR$^{40,45-47}$ and the Central Indian Ridge$^{21}$.

**Comparisons with TAG hydrothermal field**

The TAG hydrothermal field located at 26° 08’N on the MAR is one of the best-studied hydrothermal systems to date$^{18}$. High-temperature hydrothermal activity is currently focused at the active TAG mound$^{17}$. Similar as Longqi-1 hydrothermal field, the active TAG vent is located in the hanging wall at the inside corner (Fig. 2c) and is interpreted to be controlled by the permeable detachment fault$^{10}$. The TAG mound is located at the boundary of the extensional zone in the north and the oblique-fissured areas in the south$^{18}$. Multiple phases of extensional faulting and oblique fracture occurred within the hydrothermal field$^{17}$. Canales et al. (2007)$^{49}$ proposed that the hydrothermal field is located in a high-porosity zone in the hanging wall, where episodic movement along the fault surfaces may lead to the increased permeability of the hanging wall, periodically reactivating the hydrothermal system and contributing to its longevity$^{48}$.

We estimated stress and shear stain field of the TAG hydrothermal field through numerical approaches (Fig. 2c, d). The method is described before and the parameters can be seen in Table.1 in the supplemental material. We do not find focal mechanisms available within the TAG segment on the GCMT. Similar with Longqi segment, our results show a strong mechanical asymmetry between inside and outside corner at TAG segment. Stress rotation occurs at inside corner, deviating from spreading direction by 40°. The active TAG hydrothermal vent is located at the high-strain area, which likely enhances the regional permeability and facilitates hydrothermal circulation at this site. We find the shear strain at TAG segment seems to be more widely spread than that at Longqi segment, both at inside and outside corner.
We infer the difference is likely associated with the geometry and stability of NTD, which should be controlled by local magmatism and spreading rates.

**Implications for permeability and hydrothermal activity in the hanging wall**

Hydrothermal circulation requires the confluence of a sufficient heat source and adequate permeability\(^{50}\). On a global scale, the melt supply to the spreading center is roughly proportional to the spreading rate. However, even in relative terms, the regional variations of the crustal permeability are poorly unconstrained\(^{51}\). It has been suggested that on slow and ultraslow spreading centers where the magma supply is low, seafloor spreading is dominated by tectonic stretching. The impact of the variations of regional-scale crustal permeability on the hydrothermal circulation could be stronger than heat source\(^{52}\).

This study demonstrates a high-shear-strain area at the inside corner (Fig. 2). A recent study\(^{12}\) shows that the heat source in our study area is probably located at the depth of 13 ± 2 km below the seafloor, and that detachment fault provides a suitable pathway for hydrothermal circulation. The lithosphere-scale interpretation, however, could not fully account for the location of the active Longqi-1 hydrothermal vent and inactive Longqi-3 hydrothermal field, because both of which are located close to the termination of the detachment fault, where the permeability is reasoned to be at a maximum. Our results show that the active Longqi-1 hydrothermal vent as well as TAG hydrothermal vent are located at the area with maximal shear strain. We therefore infer that the location of the hydrothermal vent should simultaneously be affected by the regional permeability of the shallow crust which is likely enhanced due to shear strain. Detachment surface is thought to be at a depth of ~1 km beneath the Longqi-1 hydrothermal vent based on seismic observations\(^{23,25}\). The intersection of the oblique faults with detachment fault at depth, coupled with a suitable heat supply extracted by detachment fault, might control the location of the Longqi-1 hydrothermal field (Fig. 4). In the shallow crust, the high-temperature fluids rise vertically due to buoyancy along the permeable oblique faults to the seafloor (Fig. 4). The different stages of hydrothermal activity including active Longqi-1 hydrothermal vent and inactive Longqi-3 hydrothermal field indicate that the location of up-flow zones changes over time in response to the movement of the oblique faults in the specific areas.

**Model limitations**

It is important to note that our numerical experiments were performed on a small spatial scale with a high resolution, an isotropic layered media and a relatively simple geometry, aiming to model the present-day stress state within a ridge-offset intersection on the MORs. Due to the lack of a depth-dependent rheology and more realistic structure of the faults, our model is limited to completely predict the depth of the shear strain beneath hydrothermal field, which reflects mechanical properties of the lithosphere and the interplay between magmatic and tectonic process. In addition, based on the isotropic media assumption, the horizontal variations of rock properties (e.g. density, cohesion and porosity) are neglected in our model, leading to it is difficult to interpret the variation of the magnitude of stresses. However, combined with field observations, we believe that our model has revealed some key first-order features of the ridge-
offset intersection on the MORs and their mechanical effects on the hydrothermal activity, which might be a useful reference model for future studies.

**Conclusions**

Combined with bathymetric data and geological observations, our model provides a potential mechanical mechanism for the hydrothermal activity at the inside corner. The main conclusions of this work are as follows:

1. Our model shows that the inside corner undergoes regional stress rotation and shear strain. Stress direction is likely deflected by more than 45°.
2. Shear strain at inside corner, coupled with a suitable heat supply, likely enhances the regional permeability, facilitates hydrothermal circulation and controls the location of Longqi-1 hydrothermal vents.

**Methods**

To investigate the present-day stress state, we developed a numerical model based on conservation equations using COMSOL Multiphysics 5.4\textsuperscript{53,54}. Assuming that accelerations other than gravity can be neglected\textsuperscript{20}, we solve the force-balance Eq. 5\textsuperscript{5} for absolute stress (no pore pressure):

\[
\nabla \cdot \sigma + F_v = 0, \quad (1)
\]

where \(\sigma\) is stress tensor, and \(F_v\) is body force. The compressive stress here is reckoned to be positive and the tensile stress is negative. We adopt an explicit hybrid finite element approach to solve force-balance equation. According to Hook’s Law, the relationship between the stress tensor and strain, \(\epsilon_{ij}\), can be described as following:

\[
\sigma_{ij} = E \epsilon_{ij}, \quad (2)
\]

Where \(i\) and \(j\) are coordinate indices (\(x, y, z\)), and \(E\) is the Young’s modulus. The shear strain components can be defined as follows:

\[
\epsilon_{ij} = \frac{1}{2} \left( \frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right), \quad (3)
\]

Where \(u\) is the displacement.

The initial numerical models comprise a 52-km-long, 53-km-wide and \(~ 4.5\)-km-deep domain and a 13-km-long, 13-km-wide and \(~ 4.5\)-km-deep domain, respectively (see in the supplementary materials). The geometry of the top surface in our model is consistent with the seafloor topography at a resolution of \(~ 50\) m. For simplification, the other surfaces are simulated as planes. The horizontal \(x\) axis is parallel to
the ridge axis, the horizontal y axis is perpendicular to the ridge axis, and the vertical z axis is directed downward. In order to eliminate boundary effects, the area of interest (seafloor in our model) is set at a significant distance away from the bottom surface. The minimum and maximum thickness of our models are 2.5 and ~ 4.5 km, respectively.

Our model contains two different stratigraphic layers (Layer A and Layer B). Based on the physical properties model developed by Hyndman and Drury (1976) and active-source seismic experiments. The density of the two layers is measured from the samples. Yang's modules and Poisson's ratio are derived from active-source P-wave velocity data based on empirical equations. The half-spreading rates are loaded on the model edges. Zero-displacement boundary condition is imposed along the x and z direction while along the y direction, both edges are slip-free. The normal stress equals the overburden pressure and shear stress is set to be zero on the bottom surface. The detail parameters of our model are shown in Table 1 in supplementary materials. Hydrostatic pressure (P) is applied to the top surface, as follows:

\[ P_{i,j} = \rho_w g z_{i,j}, \]  

where P is hydrostatic pressure, \( \rho_w \) is the water density, and z is the depth. Acceleration due to gravity (g) is 9.8 m \( \cdot \) s\(^{-2}\).

**Data Availability**

The focal mechanism data used in this study are supplied in the Supplementary Tables. The bathymetry data of TAG available on GEOMAP app. Bathymetric data of Longqi can be provided upon request by sending an e-mail to the corresponding author (taochunhuimail@163.com).

**Declarations**

**Acknowledgements**

This work was supported by, National Natural Science Foundation of China (42127807, 41906071), Scientific Research Fund of the Second Institute of Oceanography, MNR, grand no. SZ2201, Talent Cultivation Project of Zhejiang Association for Science and Technology (SKX201901) and Oceanic Interdisciplinary Program of Shanghai Jiao Tong University (SL2020MS033).

**Author contributions**

M. C. wrote down the paper. He was principal investigator of the scientific program, codesigned the experiments and set down the conceptual model. C.T. conceived of the project and contributed the interpretation and writing. Y.L. H.W. and W.Y. participated in the interpretation of the results. K.C. and Q. L. participated in setting up model. M.X. and L.L. participated in data collection. D.C. participated in geology interpretation. All authors participated in the review of this manuscript.
Competing interests

The authors declare no competing interests.

References


**Tables**

Tables 1 and 2 are available in the Supplementary Files section.

**Figures**

**Figure 1**

Legend not included with this version.
Figure 2

Legend not included with this version.
Figure 3

Legend not included with this version.
Figure 4

Legend not included with this version.

Supplementary Files

This is a list of supplementary files associated with this preprint. Click to download.
• SupplementaryMaterials.pdf